



#### **Rotational Spectroscopy (from ERA):**



#### **Rotational Spectroscopy:**



ASTR 5340 - Introduction to Radio Astronomy Contact: sscibell@nrao.edu



Energy

## Rotational Lines at Radio Wavelengths: The Best Probe of *Complex* Molecules



## Rotational Lines at Radio Wavelengths: The Best Probe of Complex Molecules





## Rotational Lines at Radio Wavelengths: The Best Probe of *Complex* Molecules



NRAC

Observatory

#### **Einstein Coefficients**

To understand how emission,  $j_v$ , and absorption,  $k_v$ , coefficients are related to the atoms and molecules in a cloud, we must introduce the Einstein coefficients!



Fig. 7.5 (ERA)



#### **Einstein Coefficients**

To understand how emission,  $j_{v}$ , and absorption,  $k_{v}$ , coefficients are related to the atoms and molecules in a cloud, we must introduce the Einstein coefficients!

 $B_{LU}$  $B_{\rm UL}$  $A_{UL}$  $E_{\mathrm{L}}$ 

 $A_{III}$  = spontaneous emission coefficient from level u to l, the probability per unit time that an atom or molecule in the upper state will emit a photon, and transition to the lower state  $(s^{-1})$ .

 $B_{III}$  = induced (or stimulated) emission coefficient from level u to l, the probability per unit time per unit intensity at a given frequency that an atom or molecule in the upper state will emit a photon and transition to the lower state ( $s^{-1}$  or  $s^{-1}$  erg<sup>-1</sup> cm<sup>2</sup> rad<sup>2</sup>).

 $B_{III}$  = absorption coefficient from level *I* to *u*; induced emission from level *u* to *I*; the probability per unit time per unit intensity at a given frequency that an atom or molecule in the lower state will absorb a photon and transition to the upper state (s<sup>-1</sup> or s<sup>-1</sup> erg<sup>-1</sup> cm<sup>2</sup> rad<sup>2</sup>).

ASTR 5340 - Introduction to Radio Astronomy Contact: sscibell@nrao.edu





Fig. 7.5 (ERA)

The photon emitted or absorbed during a transition between the upper and lower states will have energy and frequency related to,

$$E = E_{\rm U} - E_{\rm L} \tag{7.38}$$

The rate (s<sup>-1</sup>) for this process is proportional to the **profileweighted mean radiation energy density** of the surrounding radiation field,

$$\bar{u} \equiv \int_0^\infty u_\nu(\nu) \ \phi(\nu) \ d\nu \qquad (7.39)$$

Where  $\varphi(v)$  is your narrow line profile. Here a system in the lower energy stage may **absorb** a photon of frequency roughly equal to the rest frequency,  $v \approx v_0$ , and **transition to the upper state**.



$$\rightarrow \qquad B_{\rm LU}\bar{u} \qquad (7.40)$$

ASTR 5340 - Introduction to Radio Astronomy Contact: sscibell@nrao.edu



Fig. 7.5 (ERA)

The photon emitted or absorbed during a transition between the upper and lower states will have energy and frequency related to,

$$E = E_{\rm U} - E_{\rm L} \tag{7.38}$$

The rate (s<sup>-1</sup>) for this process is proportional to the **profileweighted mean radiation energy density** of the surrounding radiation field,

$$\bar{u} \equiv \int_0^\infty u_\nu(\nu) \ \phi(\nu) \ d\nu \qquad (7.39)$$

Where  $\varphi(v)$  is your narrow line profile. Here a system in the lower energy stage may **absorb** a photon of frequency roughly equal to the rest frequency,  $v \approx v_0$ , and **transition to the upper state**.



Fig. 7.5 (ERA)

 $\rightarrow \qquad B_{\rm LU}\bar{u} \qquad (7.40)$ 

Analogous for 'stimulated emission' or **negative absorption**:

$$B_{\rm UL}\bar{u}$$
 (7.41)



We can write this coefficients now in useful ways...

The first to **balance** the system, considering full thermodynamic equilibrium (TE) for  $n_{U}$  or  $n_{L}$  atoms or molecules per unit volume in 'upper' or 'lower' energy states:

$$n_{\rm U}A_{\rm UL} + n_{\rm U}B_{\rm UL}\bar{u} = n_{\rm L}B_{\rm LU}\bar{u}. \tag{7.42}$$





ASTR 5340 - Introduction to Radio Astronomy Contact: sscibell@nrao.edu Fig. 7.5 (ERA)

We can write this coefficients now in useful ways...

The first to **balance** the system, considering full thermodynamic equilibrium (TE) for  $n_{U}$  or  $n_{L}$  atoms or molecules per unit volume in 'upper' or 'lower' energy states:

$$n_{\rm U}A_{\rm UL} + n_{\rm U}B_{\rm UL}\bar{u} = n_{\rm L}B_{\rm LU}\bar{u}. \tag{7.42}$$

And the ratio of  $n_{\rm U}$  to  $n_{\rm L}$  is fixed by the **Boltzmann equation** 

$$\swarrow \frac{n_{\rm U}}{n_{\rm L}} = \frac{g_{\rm U}}{g_{\rm L}} \exp\left[-\frac{(E_{\rm U} - E_{\rm L})}{kT}\right] = \frac{g_{\rm U}}{g_{\rm L}} \exp\left(-\frac{h\nu_0}{kT}\right), \quad (7.43)$$

where  $g_U$  and  $g_L$  are the numbers of distinct physical states or (statistical weights) having energies  $E_U$  and  $E_L$ 





We can write this coefficients now in useful ways...

The first to **balance** the system, considering full thermodynamic equilibrium (TE) for  $n_U'$  or  $n_L'$  atoms or molecules per unit volume in 'upper' or 'lower' energy states:

$$\left| n_{\rm U}A_{\rm UL} + n_{\rm U}B_{\rm UL}\bar{u} = n_{\rm L}B_{\rm LU}\bar{u}. \right| \qquad (7.42)$$

And the ratio of  $n_U$  to  $n_L$  is fixed by the **Boltzmann equation** 

$$\swarrow \frac{n_{\rm U}}{n_{\rm L}} = \frac{g_{\rm U}}{g_{\rm L}} \exp\left[-\frac{(E_{\rm U} - E_{\rm L})}{kT}\right] = \frac{g_{\rm U}}{g_{\rm L}} \exp\left(-\frac{h\nu_0}{kT}\right), \quad (7.43)$$

where  $g_U$  and  $g_L$  are the numbers of distinct physical states or (statistical weights) having energies  $E_U$  and  $E_L$ 



ASTR 5340 - Introduction to Radio Astronomy Contact: sscibell@nrao.edu



Fig. 7.5 (ERA)

We can write this coefficients now in useful ways...

The first to **balance** the system, considering full thermodynamic equilibrium (TE) for  $n_{U}$  or  $n_{L}$  atoms or molecules per unit volume in 'upper' or 'lower' energy states:

$$n_{\rm U}A_{\rm UL} + n_{\rm U}B_{\rm UL}\bar{u} = n_{\rm L}B_{\rm LU}\bar{u}. \qquad (7.42)$$

And the ratio of  $n_U$  to  $n_L$  is fixed by the Boltzmann equation

$$\sum_{n_{\rm L}} \frac{n_{\rm U}}{n_{\rm L}} = \frac{g_{\rm U}}{g_{\rm L}} \exp\left[-\frac{(E_{\rm U} - E_{\rm L})}{kT}\right] = \frac{g_{\rm U}}{g_{\rm L}} \exp\left(-\frac{h\nu_0}{kT}\right), \quad (7.43)$$

where  $g_U$  and  $g_L$  are the numbers of distinct physical states or (statistical weights) having energies  $E_U$  and  $E_L$ 



**Excitation Temperature** is defined by the **Boltzmann equation** and gives the ratio of the populations in each level  $(T_{ex} \text{ or } T_X)$ 

$$T_{ex} = rac{h
u/k}{\lnrac{n_l g_u}{n_u g_l}},$$

Fig. 7.5 (ERA)



Remember our radiative transfer equation:

$$rac{dI_{
u}}{ds} = -\kappa I_{
u} + j_{
u},$$
 (2.27 and 7.52)

We can construct emission and absorption coefficients from the Einstein coefficients,

$$\kappa = \left(\frac{h\nu_0}{c}\right) \left(n_{\rm L}B_{\rm LU} - n_{\rm U}B_{\rm UL}\right)\phi\left(\nu\right). \tag{7.55}$$

$$\frac{dI_{\nu}}{ds} = j_{\nu} = \left(\frac{h\nu_0}{4\pi}\right) n_{\rm U}A_{\rm UL}\phi\left(\nu\right). \tag{7.56}$$

And write the line radiative transfer equation,

$$\bigwedge \frac{dI_{\nu}}{ds} = -\left(\frac{h\nu_0}{c}\right)\left(n_{\rm L}B_{\rm LU} - n_{\rm U}B_{\rm UL}\right)\phi\left(\nu\right)I_{\nu} + \left(\frac{h\nu_0}{4\pi}\right)n_{\rm U}A_{\rm UL}\phi\left(\nu\right).$$
(7.47)



Remember our radiative transfer equation:

$$rac{dI_{
u}}{ds} = -\kappa I_{
u} + j_{
u},$$
 (2.27 and 7.52)

We can construct emission and absorption coefficients from the Einstein coefficients,

$$\kappa = \left(\frac{h\nu_0}{c}\right) \left(n_{\rm L}B_{\rm LU} - n_{\rm U}B_{\rm UL}\right)\phi\left(\nu\right). \tag{7.55}$$

$$\frac{dI_{\nu}}{ds} = j_{\nu} = \left(\frac{h\nu_0}{4\pi}\right) n_{\rm U} A_{\rm UL} \phi\left(\nu\right). \tag{7.56}$$

And write the line radiative transfer equation,

-

$$\int \frac{dI_{\nu}}{ds} = -\left(\frac{h\nu_0}{c}\right)\left(n_{\rm L}B_{\rm LU} - n_{\rm U}B_{\rm UL}\right)\phi\left(\nu\right)I_{\nu} + \left(\frac{h\nu_0}{4\pi}\right)n_{\rm U}A_{\rm UL}\phi\left(\nu\right).$$
(7.47)

How does this determine what we observe with our radio telescope?



Remember our radiative transfer equation:

$$rac{dI_{
u}}{ds} = -\kappa I_{
u} + j_{
u},$$
 (2.27 and 7.52)

We can simplify by defining a source function, S<sub>v</sub>, and writing in terms of the **optical depth**,  $\tau_v$ 

$$rac{dI_{
m v}}{d au_{
m v}} = I_{
m v} - S_{
m v}$$
 where  $S_{
m v} = rac{j_{
m v}}{k_{
m v}}$  and  $au_{
m v} = \int\limits_{
m o}^{L} k_{
m v} dz$ 



Remember our radiative transfer equation:

$$rac{dI_{
u}}{ds} = -\kappa I_{
u} + j_{
u},$$
 (2.27 and 7.52)

We can simplify by defining a source function, S<sub>v</sub>, and writing in terms of the **optical depth**,  $\tau_v$ 

$$\frac{dI_{v}}{d\tau_{v}} = I_{v} - S_{v}$$
 where  $S_{v} = \frac{j_{v}}{k_{v}}$  and  $\tau_{v} = \int_{o}^{L} k_{v} dz$ 

by multiplying both sides by  $e^{\tau v}$  a general solution to the equation of radiative transfer is found,

$$I_{\nu} = I_{\nu}(0)e^{-\tau_{\nu}} + \int_{0}^{\tau} \exp[-(\tau_{\nu} - \tau')]S_{\nu}d\tau'$$



Remember our radiative transfer equation:

$$\frac{dI_{\nu}}{ds} = -\kappa I_{\nu} + j_{\nu},$$
 (2.27 and 7.52)

We can simplify by defining a source function, S<sub>v</sub>, and writing in terms of the **optical depth**,  $\tau_v$ 

$$rac{dI_{
m v}}{d au_{
m v}} = I_{
m v} - S_{
m v}$$
 where  $S_{
m v} = rac{j_{
m v}}{k_{
m v}}$  and  $au_{
m v} = \int\limits_{
m o}^{
m L} k_{
m v} dz$ 

by multiplying both sides by  $e^{\tau v}$  a general solution to the equation of radiative transfer is found,

$$I_{\nu} = I_{\nu}(0)e^{-\tau_{\nu}} + \int_{0}^{\tau} \exp[-(\tau_{\nu} - \tau')]S_{\nu}d\tau'$$

The sum of the background intensity attenuated by the interstellar medium plus the integrated emission attenuated by the effective absorption due to the ISM between the point of emission and the observer



#### **Radiative Transfer with Einstein Coefficients**

A standard observing mode in astronomy is to measure the intensity along a line-of-sight, LOS, to the object of interest (e.g., an interstellar cloud) and then along a LOS just off the object,

$$\Delta I_{v} = I_{v}^{ON} - I_{v}^{OFF} = [S_{v} - I_{v}](1 - e^{-\tau_{v}})$$





#### **Radiative Transfer with Einstein Coefficients**

A standard observing mode in astronomy is to measure the intensity along a line-of-sight, LOS, to the object of interest (e.g., an interstellar cloud) and then along a LOS just off the object,

$$\Delta I_{v} = I_{v}^{ON} - I_{v}^{OFF} = [S_{v} - I_{v}](1 - e^{-\tau_{v}})$$

Now, in terms of Einstein coefficients,

$$S_{v} = \frac{j_{v}}{k_{v}} = \frac{n_{u}A_{ul}}{n_{l}B_{lu} - n_{u}B_{ul}} = \frac{A_{ul}/B_{ul}}{(n_{l}B_{lu}/n_{u}B_{ul}) - 1}$$

And, assuming the gas is in TE, the **Boltzmann** equation and equations of detailed balance can be substituted to yield,

$$S_{\nu} = \frac{(2h\nu^3/c^2)}{g_u n_l/g_l n_u - 1} \Leftrightarrow \frac{2h\nu^3}{c^2} \frac{1}{e^{h\nu/kT_{ex}} - 1}$$

ASTR 5340 - Introduction to Radio Astronomy Contact: sscibell@nrao.edu



Which is?!



#### **Radiative Transfer with Einstein Coefficients**

A standard observing mode in astronomy is to measure the intensity along a line-of-sight, LOS, to the object of interest (e.g., an interstellar cloud) and then along a LOS just off the object,

$$\Delta I_{v} = I_{v}^{ON} - I_{v}^{OFF} = [S_{v} - I_{v}](1 - e^{-\tau_{v}})$$

Now, in terms of Einstein coefficients,

$$S_{v} = \frac{j_{v}}{k_{v}} = \frac{n_{u}A_{ul}}{n_{l}B_{lu} - n_{u}B_{ul}} = \frac{A_{ul}/B_{ul}}{(n_{l}B_{lu}/n_{u}B_{ul}) - 1}$$

And, assuming the gas is in TE, the **Boltzmann** equation and equations of detailed balance can be substituted to yield,

$$S_{\nu} = \frac{(2h\nu^3/c^2)}{g_u n_l/g_l n_u - 1} \Leftrightarrow \frac{2h\nu^3}{c^2} \frac{1}{e^{h\nu/kT_{ex}} - 1}$$







#### **Radiative Transfer with Einstein Coefficients**

In the case of **"local thermodynamic** equilibrium" (LTE) the radiative field is defined as a blackbody (Planck Function) with a constant excitation temperature,  $T_{ex}$ (where  $T_k \sim T_{ex}$ )

$$\Delta I_{v} = [B_{v}(T_{ex}) - B_{v}(T_{BG})](1 - e^{-\tau_{v}})$$





#### **Radiative Transfer with Einstein Coefficients**

In the case of **"local thermodynamic** equilibrium" (LTE) the radiative field is defined as a blackbody (Planck Function) with a constant excitation temperature,  $T_{ex}$ (where  $T_k \sim T_{ex}$ )

$$\Delta I_{v} = [B_{v}(T_{ex}) - B_{v}(T_{BG})](1 - e^{-\tau_{v}})$$

In the Rayleigh–Jeans limit:

$$B_{\nu}(T)=\frac{2\nu^2kT}{c^2}\propto T$$

$$\Delta T = [T_{ex} - T_{BG}](1 - e^{-\tau_v})$$





#### **Radiative Transfer with Einstein Coefficients**

In the case of **"local thermodynamic** equilibrium" (LTE) the radiative field is defined as a blackbody (Planck Function) with a constant excitation temperature,  $T_{ex}$ (where  $T_k \sim T_{ex}$ )

$$\Delta I_{v} = [B_{v}(T_{ex}) - B_{v}(T_{BG})](1 - e^{-\tau_{v}})$$





#### **Radiative Transfer with Einstein Coefficients**

In the case of **"local thermodynamic** equilibrium" (LTE) the radiative field is defined as a blackbody (Planck Function) with a constant excitation temperature,  $T_{ex}$ (where  $T_k \sim T_{ex}$ )

$$\Delta I_{v} = [B_{v}(T_{ex}) - B_{v}(T_{BG})](1 - e^{-\tau_{v}})$$

In the limiting cases,

 $\Delta I_{v} = [B_{v}(T_{ex}) - B_{v}(T_{BG})] \quad \text{for } \tau_{v} \gg 1$  $\Delta I_{v} = [B_{v}(T_{ex}) - B_{v}(T_{BG})]\tau_{v} \quad \text{for } \tau_{v} \ll 1$ 



#### **Radiative Transfer with Einstein Coefficients**

In the case of **"local thermodynamic** equilibrium" (LTE) the radiative field is defined as a blackbody (Planck Function) with a constant excitation temperature,  $T_{ex}$ (where  $T_k \sim T_{ex}$ )

$$\Delta I_{v} = [B_{v}(T_{ex}) - B_{v}(T_{BG})](1 - e^{-\tau_{v}})$$

In the limiting cases,

$$\Delta I_{\nu} = [B_{\nu}(T_{ex}) - B_{\nu}(T_{BG})] \quad \text{for } \tau_{\nu} \gg 1$$
$$\Delta I_{\nu} = [B_{\nu}(T_{ex}) - B_{\nu}(T_{BG})]\tau_{\nu} \quad \text{for } \tau_{\nu} \ll 1$$

**Optically thick lines** (e.g., <sup>12</sup>CO) or dust emission can be used to determine the **temperature** of a cloud!

**Optically thin lines** (e.g., <sup>13</sup>CO or C<sup>18</sup>O) or dust emission is directly proportional to the **cloud's optical depth** (and thus column density and cloud mass)!



$$\Delta T_{\rm b} \approx (T_{\rm x} - T_{\rm c}) \tau_0 = \left(\frac{T_{\rm x} - T_{\rm c}}{T_{\rm x}}\right) N_{\rm L} \frac{(\ln 2)^{1/2}}{4\pi^{3/2}} \frac{hc^3}{k\nu_0^2} \frac{g_{\rm U}}{g_{\rm L}} \frac{A_{\rm UL}}{\Delta \nu}$$
(7.139)  
**Optically thick lines** (e.g., <sup>12</sup>CO) or dust emission can be used to determine the **temperature** of a cloud!  
**Optically thin lines** (e.g., <sup>13</sup>CO or C<sup>18</sup>O) or dust emission is directly proportional to the **cloud's optical depth** (and thus column density and cloud mass)!



Radiative Transfer with Einstein Coefficients

# $\Delta T_{\rm b} \approx (T_{\rm x} - T_{\rm c}) \tau_0 = \left(\frac{T_{\rm x} - T_{\rm c}}{T_{\rm x}}\right) N_{\rm L} \frac{(\ln 2)^{1/2}}{4\pi^{3/2}} \frac{hc^3}{k\nu_0^2} \frac{g_{\rm U}}{g_{\rm L}} \frac{A_{\rm UL}}{\Delta v}$ (7.139)

To get column density,  $N_{i}$ , you need to know the physical states of your molecule and calculate the spontaneous emission <u>coefficient</u>

**Optically thick lines** (e.g., <sup>12</sup>CO) or dust emission can be used to determine the **temperature** of a cloud!

**Optically thin lines** (e.g., <sup>13</sup>CO or C<sup>18</sup>O) or dust emission is directly proportional to the cloud's optical depth (and thus column density and cloud mass)!



The spontaneous emission coefficient be calculated by writing in terms of the **mean electric dipole moment**,

$$A_{\rm UL} = \frac{64\pi^4}{3hc^3} |\mu_{\rm UL}|^2 \nu^3.$$
 (7.131)

For a linear rotating molecule,

$$|\mu_{J\to J-1}|^2 = \frac{\mu^2 J}{2J+1}.$$
 (7.132)

$$\left(\frac{A_{J\to J-1}}{\mathrm{s}^{-1}}\right) \approx 1.165 \times 10^{-11} \left|\frac{\mu}{\mathrm{D}}\right|^2 \left(\frac{J}{2J+1}\right) \left(\frac{\nu}{\mathrm{GHz}}\right)^3. \quad (7.133)$$



#### **Radiative Transfer with Einstein Coefficients**

The spontaneous emission coefficient be calculated by writing in terms of the **mean electric dipole moment**,

$$A_{\rm UL} = \frac{64\pi^4}{3hc^3} |\mu_{\rm UL}|^2 \nu^3.$$
 (7.131)

For a linear rotating molecule,

$$|\mu_{J\to J-1}|^2 = \frac{\mu^2 J}{2J+1}.$$
 (7.132)

$$\left(\frac{A_{J\to J-1}}{s^{-1}}\right) \approx 1.165 \times 10^{-11} \left|\frac{\mu}{D}\right|^2 \left(\frac{J}{2J+1}\right) \left(\frac{\nu}{GHz}\right)^3$$
. (7.133)



Example for CO (1-0):

$$A_{10} \approx 1.165 \times 10^{-11} \cdot 0.11^2 \cdot \left(\frac{1}{3}\right) \cdot 115^3$$
  
 $\approx 7.1 \times 10^{-8} \text{ s}^{-1}$ 

Solving for **critical density** assumes a typical collisional cross section and average velocity for  $H_2$  molecules,

$$n^* \approx \frac{A_{\rm UL}}{\sigma v} \approx \frac{7 \times 10^{-8} \text{ s}^{-1}}{10^{-15} \text{ cm}^2 \cdot 5 \times 10^4 \text{ cm s}^{-1}}$$
  
 $\approx 1.4 \times 10^3 \text{ cm}^{-3}.$ 



#### **Radiative Transfer with Einstein Coefficients**

Related to  $A_{UL}$ , the **upper energies** give clues into what type of environments are molecular lines are likely to emit at and are directly connected to the minimum gas temperature needed for significant collisional excitation,

$$T_{\min} \sim \frac{E_{\text{rot}}}{k}$$

$$\sim \frac{J(J+1)h^2}{2 \cdot 4\pi^2 Ik} = \frac{hJ}{4\pi^2 I} \frac{h(J+1)}{2k} = \frac{E_U}{k}$$
(7.116 & 7.118)



ASTR 5340 - Introduction to Radio Astronomy Contact: sscibell@nrao.edu



Fig. 7.15 (ERA)

#### **Radiative Transfer with Einstein Coefficients**

$$\Delta T_{\rm b} \approx (T_{\rm x} - T_{\rm c}) \tau_0 = \left(\frac{T_{\rm x} - T_{\rm c}}{T_{\rm x}}\right) N_{\rm L} \frac{(\ln 2)^{1/2}}{4\pi^{3/2}} \frac{hc^3}{k\nu_0^2} \frac{g_{\rm U}}{g_{\rm L}} \frac{A_{\rm UL}}{\Delta \nu}$$
(7.139)  
$$T_{\rm min} \sim \frac{E_{\rm rot}}{k}$$
$$\sim \frac{J \left(J + 1\right) h^2}{2 \cdot 4\pi^2 I k} = \frac{hJ}{4\pi^2 I} \frac{h \left(J + 1\right)}{2k} = \frac{E_{\rm U}}{k}$$

(7.116 & 7.118)

To get column density, N<sub>L</sub>, you need to know the <u>physical states</u> of your molecule and <u>calculate the</u> <u>spontaneous emission</u> <u>coefficient</u>

In practice... we look up these terms in Splatalogue! <u>https://splatalogue.online</u>





Fig. 7.16 (ERA)

What spectral line are we looking at? Use Splatalogue! <u>https://splatalogue.online</u>



















